

19
ИНСТИТУТ ЯДЕРНОЙ ФИЗИКИ
СО АН СССР

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WEAK NEUTRAL CURRENTS OF NEW QUARKS

IN e^+e^- ANNIHILATION

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ПРЕПРИНТ ИЯФ 79 - 25

Новосибирск

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abstract

P-odd dependence on the longitudinal polarization of colliding particles of the total cross-section of hadron production in e^+e^- collisions near Υ^- and Υ^+ -peaks is calculated. At Υ^- -peak the expected relative magnitude of the effect constitutes 1.6%.

Obtaining of longitudinal polarization of electrons and positrons in storage rings^{/1-3/} would open great prospects for the study of the weak neutral currents structure in e^+e^- collisions. Experiments on the search for weak interactions effects in colliding electron-positron beams were often discussed previously^{/4-20/}. Here we wish to call attention to the importance of study of P - odd dependence of the total cross-section of hadron production near Υ^- and Υ^+ -peaks on longitudinal polarization of electrons and positrons. This experiment would give valuable information on weak interaction of new quarks and electrons due to neutral currents. For u- and d-quarks the information on such an interaction is obtained already in the atomic experiments^{/21/} and in deep inelastic electron scattering^{/22/}, as to c- and b-quarks the colliding beams experiments would be of unique interest indeed.

Write down the Lagrangian for the interaction of electrons and quarks with photon and neutral vector boson as

$$\mathcal{L} = e A_\mu [\bar{u}_e \gamma_\mu u_e - Q \bar{u}_q \gamma_\mu u_q] + Z_\mu [\bar{u}_e \gamma_\mu (g_V + g_A \gamma_5) u_e + \bar{u}_q \gamma_\mu (h_V + h_A \gamma_5) u_q] \quad (1)$$

Standard calculations lead to the following dependence of the total cross-section of hadron production near a peak corresponding to $q\bar{q}$ bound state on the degree of longitudinal polarization of electron ξ_- and positron ξ_+ :

$$\sigma = \sigma_0 \left[1 - \xi_- \xi_+ - \frac{2 h_V g_A m^2}{4\pi \alpha M^2} (\xi_- - \xi_+) \right] \quad (2)$$

where σ_0 is the resonance cross-section for unpolarized beam. We assume here that the resonance mass m is much smaller than that of Z-boson M . Parity violation is characterized by the

quantity

$$\eta(\bar{q}_-, \bar{q}_+) = \frac{\sigma(\bar{q}_-, \bar{q}_+) - \sigma(-\bar{q}_-, -\bar{q}_+)}{\sigma(\bar{q}_-, \bar{q}_+) + \sigma(-\bar{q}_-, -\bar{q}_+)} \quad (3)$$

In our case

$$\eta(\bar{q}_-, \bar{q}_+) = - \frac{2h\nu g_A m^2}{4\pi\alpha M^2 Q} \frac{\bar{q}_- - \bar{q}_+}{1 - \bar{q}_- \bar{q}_+} \quad (4)$$

In the Weinberg-Salam model

$$g_A = - \frac{e}{2\sin 2\theta}, \quad h\nu = \frac{e}{2\sin 2\theta} \left[\frac{Q}{|Q|} - 4Q\sin^2\theta \right] \quad (5)$$

and for the total longitudinal polarization of both or only one of the initial particles we find for $q\bar{q}$ resonance

$$\eta \equiv \eta(1, -1) = \eta(1, 0) = \eta(0, -1) = \frac{\sqrt{2} G m^2}{8\pi\alpha |Q|} (1 - 4|Q|\sin^2\theta) \quad (6)$$

At $\sin^2\theta = \frac{1}{4}$ ¹⁾ this quantity constitutes for Ψ - and Υ -peaks correspondingly

$$\eta_\Psi = \frac{\sqrt{2} G m^2}{16\pi\alpha} = 4 \cdot 10^{-4}, \quad \eta_\Upsilon = \frac{\sqrt{2} G m^2}{4\pi\alpha} = 1.6 \cdot 10^{-2} \quad (7)$$

Note that for Υ -meson the effect is not much smaller than in Ψ -peak:

$$\eta_\Upsilon = \frac{\sqrt{2} G m^2}{4\pi\alpha} = 1.8 \cdot 10^{-4} \quad (8)$$

We have assumed here that in a peak non-resonant contribution to the hadron production cross-section can be neglected.

¹⁾ It is curious that at such mixing angle the product of weak and electric charges is the same for all quarks: $-eQh\nu = -\frac{2\sqrt{3}}{24}e^2$. On the other hand, at $\sin^2\theta = \frac{3}{8}$ (this value of the mixing parameter is predicted by the unified model of strong, electromagnetic and weak interactions based on SU(5) symmetry^[23]) weak charge of the quarks with $Q = 2/3$ turns to zero, and the discussed effect in Ψ -peak vanishes as well.

Near Ψ -peak the non-resonant effect due to the contribution of u-, d- and s-quarks to the cross-section constitutes

$$\eta = \frac{\sqrt{2} G S}{4\pi\alpha} (1 - 2\sin^2\theta) \approx 8 \cdot 10^{-4} \quad (9)$$

and near Υ -peak where one should take into account c-quark as well, it is equal to

$$\eta = \frac{\sqrt{2} G S}{4\pi\alpha} \frac{9 - 20\sin^2\theta}{10} \approx 0.6 \cdot 10^{-2} \quad (10)$$

Relatively large expected value of the effect discussed in Υ -resonance and possibility to get large statistics in Υ -peak make the corresponding experiments sufficiently realistic, of course, if the longitudinal polarization will be obtained in colliding beams.

Note in conclusion that near narrow resonance of the Υ -peak type charge asymmetry of muons in the reaction $e^+e^- \rightarrow \mu^+\mu^-$ due to the interference of electromagnetic and weak amplitudes turns to be very small²⁾. As it was noted by A.P. Onuchin, this circumstance can be used effectively for control experiments at the search for the charge asymmetry outside of peaks.

²⁾ For the charge asymmetry caused by radiative corrections this fact was pointed out previously in the references^[24,25].

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Работа поступила - 2 апреля 1979 г.

Ответственный за выпуск - С.Г.ПОПОВ
 Подписано к печати 23.IV-1979 г. МН 02775
 Усл. 0,4 печ.л., 0,3 учетно-изд.л.
 Тираж 200 экз. Бесплатно
 Заказ № 25.

Отпечатано на ротационной машине ИЯФ СО АН СССР